

Crystal coherence length effects on the infrared optical response of MgO thin films

J. F. Ihlefeld,^{1,a)} J. C. Ginn,¹ D. J. Shelton,² V. Matias,³ M. A. Rodriguez,¹ P. G. Kotula,¹ J. F. Carroll III,¹ G. D. Boreman,² P. G. Clem,¹ and M. B. Sinclair¹

¹Sandia National Laboratories, Albuquerque, New Mexico 87185, USA

²College of Optics and Photonics/CREOL, University of Central Florida, Orlando, Florida 32816, USA

³Los Alamos National Laboratory, Los Alamos, New Mexico 87545, USA

(Received 29 September 2010; accepted 23 October 2010; published online 12 November 2010)

The role of crystal coherence length on the infrared optical response of MgO thin films was investigated with regard to Reststrahlen band photon-phonon coupling. Preferentially (001)-oriented sputtered and evaporated ion-beam assisted deposited thin films were prepared on silicon and annealed to vary film microstructure. Film crystalline coherence was characterized by x-ray diffraction line broadening and transmission electron microscopy. The infrared dielectric response revealed a strong dependence of dielectric resonance magnitude on crystalline coherence. Shifts to lower transverse optical phonon frequencies were observed with increased crystalline coherence. Increased optical phonon damping is attributed to increasing granularity and intergrain misorientation. © 2010 American Institute of Physics. [doi:10.1063/1.3515901]

The study of optical devices built with low-loss negative-permittivity and permittivity-near-zero components is generating great interest, particularly for metamaterial applications.^{1,2} To implement materials into designs utilizing these phenomena, knowledge of the effect of defects and microstructure on optical response is necessary; this is especially true as device dimensions approach typical defect sizes on the order of nanometers. The coupling of infrared (IR) photons with near-zero wavevector optical phonons results in a band of high-reflectivity, commonly denoted the Reststrahlen band, in partially ionic materials. In this band the real portion of permittivity crosses zero and can become strongly negative. As the transverse optical phonon oscillations responsible for this response are collective phenomena requiring a degree of lattice coherence, knowledge of the role of crystal coherence length on the photon-phonon interactions is pertinent. Surface asperities, bulk porosity, and particle size are known to affect the spectral shape and intensity of the infrared optical response of ionic crystals, ceramics, and powders.^{3–9} These correlations, however, have been predominantly investigated for relatively large-scale micron-size features. The role of submicron defects including dislocations and crystallographic disorder are well known to affect phonon lifetimes and scattering (see, for example, Ref. 10 and citations within) but have been minimally studied in the context of the Reststrahlen band response in thin films.¹¹

Magnesium oxide (MgO) crystallizes in the rocksalt structure (*Fm3m*) and, owing to its high symmetry and two ion basis, possesses relatively simple phonon dispersion and Reststrahlen spectra. The Reststrahlen band in MgO occurs for incident optical wavelengths between 13.6 μm (0.091 eV) and 25 μm (0.049 eV).¹² The effect of defects on phonon damping and IR optical response has been evaluated for ion implanted MgO single crystals previously.¹³ In this report we present the effect of crystal coherence length on the IR optical response of dense polycrystalline MgO thin films with equivalent thickness and surface roughness.

200 nm thick MgO thin films were prepared on (001)-oriented silicon substrates via 30° off-axis radio frequency (rf) magnetron sputtering and ion-beam assisted deposition (IBAD). The substrates were subjected to a buffered HF etch and de-ionized-H₂O rinse prior to loading into the load-locked sputter chamber to remove the native SiO₂ surface. A polycrystalline crystal coherence length series was prepared from films deposited by RF sputtering MgO from a single-phase sintered target in 5 mTorr of argon at room temperature with a sputter power density of 3.7 W/cm². The wafer was subsequently divided and sections were processed at 200–800 °C in 200 °C intervals in air for 1 h. Thermodynamic calculations predict phase stability of MgO with silicon in this temperature range with no chemical reactions, suggesting that abrupt interfaces with no secondary phases should form in this process.¹⁴ A 200 nm thick biaxially textured (001)-oriented film was prepared via the IBAD method as described elsewhere.¹⁵ Briefly, a 5 nm amorphous Y₂O₃ layer is deposited by evaporation followed by a ~5 nm IBAD-MgO layer. A 200 nm thick homoepitaxial layer is then deposited at 600 °C by electron-beam sublimation from a MgO source. Film roughness was assessed with atomic force microscopy (AFM). Crystal coherence length was characterized with x-ray diffraction (XRD) in both Bragg-Brentano (Philips MRD) and grazing-incidence (Siemens D-500 with 2° alpha angle) geometries and transmission electron microscopy (TEM) (FEI Tecnai F30-ST) in TEM and scanning TEM modes. Crystal coherence lengths were determined using x-ray line broadening and Scherrer's formula.¹⁶ The structure factor for MgO limits the number of diffraction peaks and therefore precluded the use of the more rigorous Williamson-Hall¹⁷ method and subsequent decoupling of strain and coherence. Infrared optical properties were measured using an infrared variable angle spectroscopic ellipsometer (Woollam IR-VASE) at 50°, 55°, and 60° incidence angles for all samples over the range of 1.5–35 μm . A least-squares fit for a Lorentzian resonator at the transverse optic phonon mode wavelength and a weak Gaussian fit in the mid-IR was performed to extract the com-

^{a)}Electronic mail: jihlefe@sandia.gov.

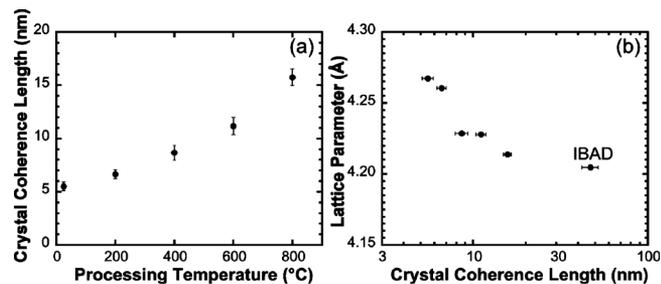


FIG. 1. (a) Crystal coherence length calculated from XRD FWHM of the 200-reflection of sputtered MgO films as-deposited and postprocessed in air from 200 °C to 800 °C. (b) Lattice parameter calculated from the 200-reflection of MgO as a function of crystal coherence length.

plex refractive index components from the ellipsometer parameters. The effect of crystal coherence length on the optical response was quantified by fitting the imaginary component of permittivity calculated from the measured optical response with complex permittivity values derived from a classical oscillator model given in Eq. (1),¹⁸

$$\varepsilon(\omega) = \varepsilon_{\infty} + \frac{\varepsilon_0 - \varepsilon_{\infty}}{\left[1 + \left(\frac{\omega^2}{\omega_{\text{TO}}^2} \right)^2 \right] - i \left(\frac{\omega\gamma}{\omega_{\text{TO}}^2} \right)}, \quad (1)$$

where ε_0 , ε_{∞} , γ , ω , and ω_{TO} are the low frequency permittivity, high frequency permittivity, damping constant, incident photon frequency, and transverse optical phonon frequency, respectively. Values of 9.86, 2.95, and 11.8 THz were used for the low frequency permittivity, high frequency permittivity, and transverse optic phonon frequency, respectively,¹⁹ as starting points and fits were introduced to determine damping constants and phonon frequencies.

Figure 1(a) shows a plot of the coherent scattering lengths for the polycrystalline MgO samples processed from room temperature to 800 °C and reveals a trend of increasing length with processing temperature. Measured coherent scattering lengths ranged from 5.5 nm (as-deposited sample) to 15.8 nm (800 °C processed sample). Identical scattering lengths were calculated using the 200 MgO x-ray reflection for both grazing-incidence and Bragg–Brentano diffraction geometries, suggesting that there is minimal directional dependence of coherent scattering lengths and that the features contributing to scattering are equiaxed. A coherent scattering length of 47 nm was measured for the (001)-oriented IBAD film (note that this is the minimum possible scattering length value as strain contributions could not be decoupled). Additionally, an omega rocking curve full width at half maximum (FWHM) value of 1.2° for the 200-reflection and 4.7° FWHM for the 220 reflection in phi were measured. MgO lattice parameters were calculated from the 200 peak position using the silicon 400 peak as a reference for sample displacement error correction²⁰ with the results shown in Fig. 1(b). A systematic reduction in lattice parameter with crystalline coherence increase is observed. These data indicate that the higher the degree of crystalline perfection, the more closely the lattice parameter approaches the accepted single crystal value of 4.21 Å. It is likely that the grain boundary disorder in the fine crystallite size samples results in expanded lattice parameters, as has been suggested previously for fine-grained MgO.⁹

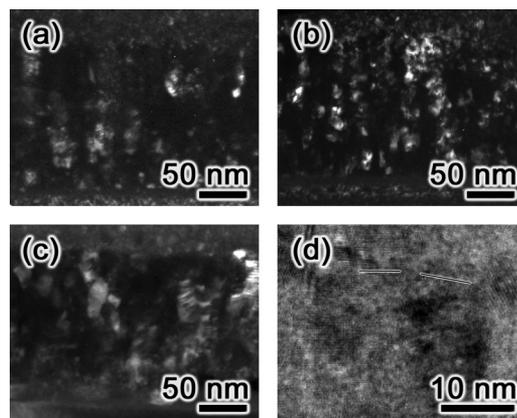


FIG. 2. Dark-field TEM images from (a) as-deposited, (b) 400 °C processed, and (c) 800 °C processed sputtered MgO thin films. (d) High-resolution TEM image of an as-deposited MgO film. Lines highlight (200) type planes for two adjacent grains.

TEM analysis of the thermally processed sputtered samples revealed that the samples were comprised of fine crystallites within larger columns consisting of preferentially crystallographic texturing with (001)-orientation. Dark-field images in Fig. 2 show bright regions where the beam satisfies diffraction conditions for individual crystallites in the as-deposited [Fig. 2(a)], 400 °C processed [Fig. 2(b)], and 800 °C processed [Fig. 2(c)] samples. Qualitatively, the average sizes of the observed individual crystallites agree well with the coherent scattering lengths measured via XRD. High-resolution images revealed small angle grain boundaries ($\sim 1^\circ - 2^\circ$) within the columnar textured regions with domain sizes nearly identical to the coherent x-ray scattering length [Fig. 2(d)]. Lines highlighting two adjacent (002)-type planes have been drawn showing small angle tilts separating adjacent crystallites. Between larger (001)-oriented columns fine-scale porosity was observed, however no obvious change in density or porosity distribution was observed with annealing condition. Additionally, no change in thickness could be identified, suggesting that film density remained constant. No difference in film roughness was observed for the sputtered and thermally processed samples. The rms roughness values measured by AFM over 9 μm^2 were 6.86 nm \pm 0.26 nm.

The real and imaginary components of the optical permittivity between 2 and 35 μm wavelengths are shown in Fig. 3 for the polycrystalline samples that were as-deposited and those processed at 400 and 800 °C in addition to the (001)-oriented IBAD-deposited sample. A single resonance

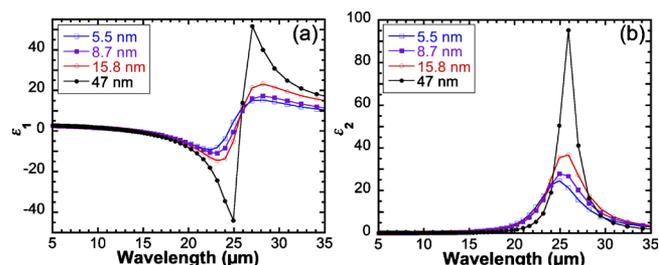


FIG. 3. (Color online) (a) Real and (b) imaginary components of the dielectric spectra obtained from infrared-spectroscopic ellipsometry analysis of thermally processed polycrystalline 200 nm thick MgO films and a 200 nm thick (001)-oriented biaxially textured IBAD film.

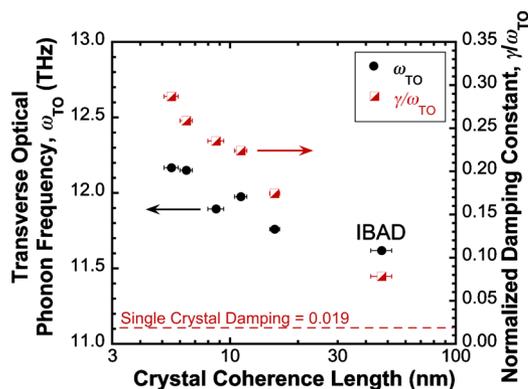


FIG. 4. (Color online) Transverse optical phonon frequency (closed circles) and normalized damping constant (open squares) for 200 nm thick MgO films as a function of crystal coherence length.

located at approximately 25 μm is observed with no evidence of additional features that could be attributed to surface modes, as are commonly observed in fine-grained powder samples of similar dimensions at shorter wavelengths than the transverse optic mode.^{3,9} The magnitudes of the polarization resonance (ϵ_1) and absorption (ϵ_2) scale with coherent crystal length, as is evident in comparing the thermally processed samples with the IBAD sample. An increase in peak resonance amplitude by greater than $4.5\times$ is observed between samples with 5.5 nm ($\epsilon_1 = -9.1$) and 47 nm ($\epsilon_1 = -44$) crystalline coherence lengths. Additionally, a shift in the transverse optic phonon mode toward longer wavelengths (lower frequency) is observed as the crystal coherence length increases. Figure 4 shows the transverse optical phonon frequency and damping constant normalized to the transverse optical phonon frequency as a function of the coherent crystal length. Similar to the lattice parameter trend, the oscillation frequency also scales with crystallinity. This was observed previously in fine-grained MgO and the increased frequency with finer crystal size was attributed to the increased lattice parameter.⁹ An increased damping constant with decreased crystal coherence is observed; this appears to be consistent with optical phonon scattering at the defects responsible for finite crystal coherence lengths. As the textured column width and porosity levels were constant for each processing condition, this suggests that the damping constants are strongly influenced by the defects responsible for crystal coherence length limitations, e.g., grain boundaries, small angle tilt boundaries, and dislocations. It is prudent to mention that an apparent shift in the longitudinal optic mode (where ϵ_1 crosses zero) to shorter wavelengths was observed as crystalline coherence length increased. Combined with the transverse optic mode shift to longer wavelengths, this would indicate that the static and high frequency dielectric constants must also scale with crystallinity through the Lyddane–Sachs–Teller relation. Unfortunately, limited instrument resolution at long wavelengths limits our ability to rigorously assess the static dielectric constant and explore this effect further.

A line indicating the accepted room temperature single crystal normalized damping constant has been plotted in Fig. 4. It is evident that as the crystal coherence length increases, bulklike values for the normalized damping constant (0.019) are approached.²¹ It is worth noting that this suggests that even the high crystalline quality of the IBAD film is not

sufficient for a bulk damping constant; substantial film thickness and lateral dimensions of high crystalline quality may be necessary to obtain a sharp resonance and subsequent high-reflectivity characteristic of the single crystalline Reststrahlen response, as has been demonstrated in other thin film systems.²² This may have implications on the feature size in both thickness and lateral dimensions for devices designed around the Reststrahlen response.

In summary, this study demonstrates that the Reststrahlen band response in MgO thin films strongly depends on the crystalline quality of the material. Films prepared with fine-scale crystallinity possess higher frequency transverse optical phonon modes and increased damping constants. The damping constant scales inversely with crystallinity resulting in an increased transverse optic phonon resonance with incident light. The reduced optical response caused by nanoscale defects may have beneficial implications for designing epsilon near-zero infrared materials, due to the flattened response through defect scattering, and may have broad implications for accurate design of IR optical components utilizing optical phonon modes in thin film materials.

The authors wish to acknowledge experimental advice and critical review from M. Lee, P. Rakich, and G. L. Brenecka. This research was supported by the Laboratory Directed Research and Development Program at Sandia National Laboratories. This work was performed, in part, at the Center for Integrated Nanotechnologies, a U.S. Department of Energy, Office of Basic Energy Sciences user facility. Sandia is a multiprogram laboratory operated by Sandia Corporation, a Lockheed Martin Co., for the U.S. Department of Energy under Contract No. DE-AC04-94AL85000. A portion of this work (LANL) was supported by the U.S. Department of Energy's Office of Electricity Delivery and Energy Reliability.

¹A. Alu, M. G. Silveirinha, A. Salandrino, and N. Engheta, *Phys. Rev. B* **75**, 155410 (2007).

²D. Korobkin, Y. A. Urzhumov, B. Neuner, C. Zorman, Z. Zhang, I. D. Mayergoyz, and G. Shvets, *Appl. Phys. A: Mater. Sci. Process.* **88**, 605 (2007).

³R. Ruppin and R. Englman, *Rep. Prog. Phys.* **33**, 149 (1970).

⁴J. R. Aronson and A. G. Emslie, *Appl. Opt.* **12**, 2573 (1973).

⁵J. W. Salisburry and A. Wald, *Icarus* **96**, 121 (1992).

⁶S. K. Andersson and C. G. Ribbing, *Phys. Rev. B* **49**, 11336 (1994).

⁷S. K. Andersson and G. A. Niklasson, *J. Phys.: Condens. Matter* **7**, 8507 (1995).

⁸J. F. Mustard and J. E. Hays, *Icarus* **125**, 145 (1997).

⁹J. T. Luxon, D. J. Montgomery, and R. Summitt, *Phys. Rev.* **188**, 1345 (1969).

¹⁰M. Kitajima, *Crit. Rev. Solid State Mater. Sci.* **22**, 275 (1997).

¹¹S. S. Ng, Z. Hassan, M. R. Hashim, and M. E. Kordesch, *Mater. Chem. Phys.* **91**, 404 (2005).

¹²A. F. Turner, L. Chang, and T. P. Martin, *Appl. Opt.* **4**, 927 (1965).

¹³H. Ogiso, M. Nakada, S. Nakano, and J. Akedo, *Nucl. Instrum. Methods Phys. Res. B* **257**, 545 (2007).

¹⁴K. J. Hubbard and D. G. Schlom, *J. Mater. Res.* **11**, 2757 (1996).

¹⁵V. Matias, J. Hanisch, E. J. Rowley, and K. Guth, *J. Mater. Res.* **24**, 125 (2009).

¹⁶P. Scherrer, *Nachr. Ges. Wiss. Goettingen, Math.-Phys. Kl.* **2**, 96 (1918).

¹⁷G. K. Williamson and W. H. Hall, *Acta Metall.* **1**, 22 (1953).

¹⁸P. Y. Yu and M. Cardona, *Fundamentals of Semiconductors: Physics and Materials Properties* (Springer, Berlin, New York, 1996).

¹⁹G. Peckham, *Proc. Phys. Soc. London* **90**, 657 (1967).

²⁰A $\cos^2 \theta / \sin \theta$ height error correction was used.

²¹J. R. Jasperse, A. Kahan, J. N. Plendl, and S. S. Mitra, *Phys. Rev.* **146**, 526 (1966).

²²K. C. Agarwal, B. Daniel, C. Klingshirn, and M. Hetterich, *Phys. Rev. B* **73**, 045211 (2006).